Tune calculation in non-scaling FFAG

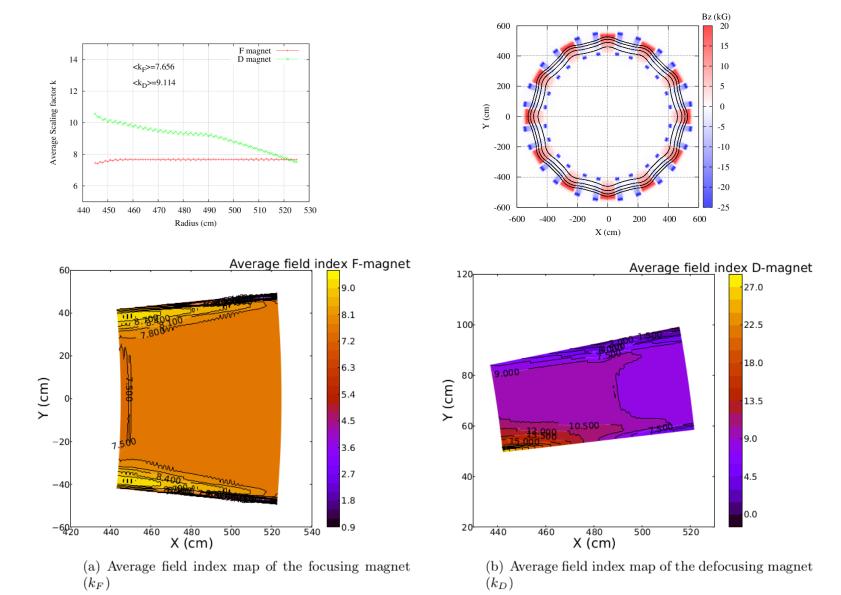
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MALEK HAJ TAHAR Doctoral student UJF-Grenoble, BNL C-AD





Scaling factors of the F and D magnet



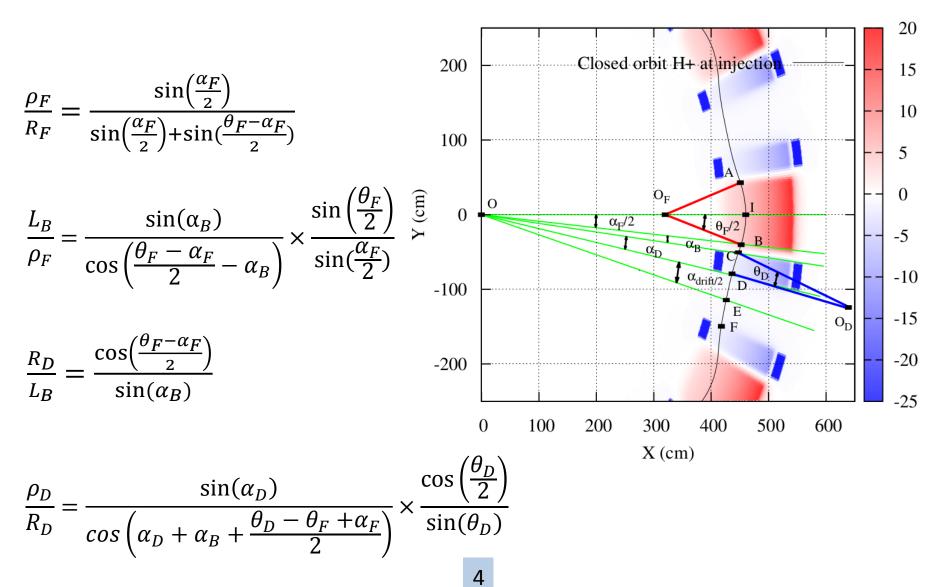
Beam stability analysis

- We use 3 different approaches to investigate the beam stability due to field errors in scaling FFAG:
- 1) Hard edge model
- 2) Bogoliubov method of averages
- 3) Zgoubi tracking model

The aim is to find a relationship between the tunes of the cell and the scaling factors k_F and k_D .

Hard edge model

Bz (kGauss)



Hard edge model

•
$$\frac{d^2x}{ds^2} + \frac{1-n}{\rho^2}x = 0$$

•
$$\frac{d^2z}{ds^2} + \frac{n}{\rho^2}z = 0$$

where
$$n(R) \approx -\frac{\rho}{B} \frac{dB}{dR} = -\frac{\rho}{R} k(R) \propto k(R)$$

$$\frac{d^2x}{ds^2} + \frac{1 + \frac{\rho}{R}k(R)}{\rho^2}x = 0$$
$$\frac{d^2z}{ds^2} - \frac{k(R)}{\rho R}z = 0$$

Conjecture:

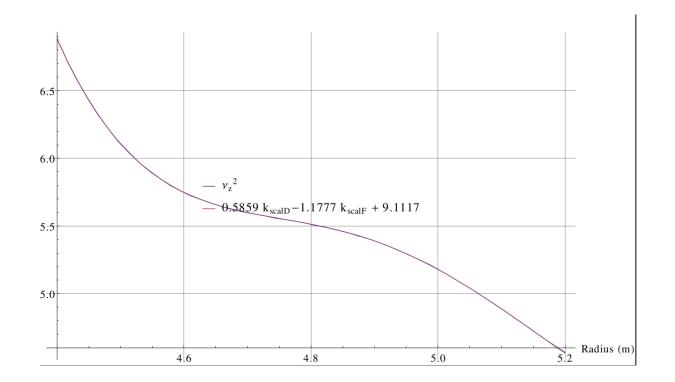
Back to Symon model:
We make the following <u>conjecture</u>:

<u>The Vertical tune</u> of a DFD cell is given by: $v_z^2 = x_1 k_F + x_2 k_D + x_3$ Analogy with $v_z^2 = -k + \frac{f^2}{2}$

<u>The horizontal tune</u> of a DFD cell is given by: $v_x^2 = x_1k_F + x_2k_D + (x_3k_F + x_4k_D + x_5)^2$

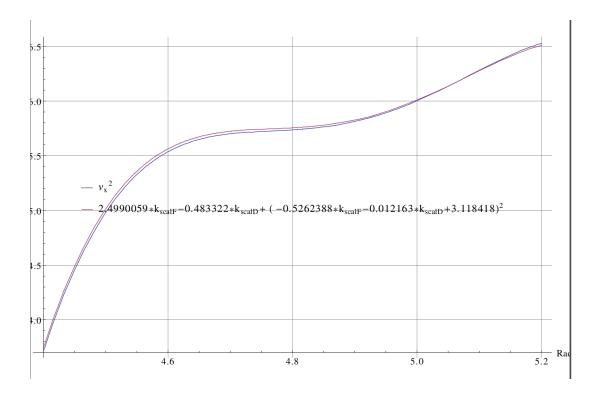
Analogy with $v_{\chi}^2 = k + 1 + A. (k + 1)^2$

Solving the Vertical Tune equation



The conjecture is valid.

Solving the Horizontal Tune equation



The conjecture is valid.

- Can the previous conjecture be proven in the general case where $\frac{\rho}{R} \neq constant$?
- This would be a generalization to the case of a non-scaling FFAG.

• For that reason, the Bogoliubov method of averages has been applied.

• Write the Hill's equation in cylindrical coordinates:

$$\begin{cases} \quad \frac{d^2x}{d\theta^2} + \left[\mu(R,\theta)^2 [1-n]\right] x = 0 \quad ; \quad \mu(R,\theta) = \frac{R}{\rho} \\ \quad \frac{d^2y}{d\theta^2} + \left[\mu(R,\theta)^2 n\right] y = 0 \end{cases}$$

$$\frac{ds}{d\theta} = \sqrt{R^2 + (\frac{dR}{d\theta})^2} \approx R$$

Arclength in cylindrical coordinates

• This can be written in the standard form:

$$\frac{d^2x}{d\theta^2} + g(R,\theta)x = 0$$

$$\nu^{2}(R) = \langle g \rangle + \frac{1}{N^{2}} \langle \tilde{g}^{2} \rangle + \dots$$
$$\approx \langle g(R,\theta) \rangle + \left\langle \left[\int \{g(R,\theta) - \langle g(R,\theta) \rangle \} d\theta \right]^{2} \right\rangle$$
$$= g_{1}(R) + g_{2}(R)$$

Calculate the field index n:

$$n = -\frac{\rho}{B} \frac{\partial B}{\partial x} = -\frac{\rho}{B} \left[\frac{\partial B}{\partial R} \frac{\partial R}{\partial x} + \frac{\partial B}{\partial \theta} \frac{\partial \theta}{\partial x} \right]$$
$$g_x(R,\theta) = \left(\frac{R}{\rho}\right)^2 \times (1-n) \approx \left(\frac{R}{\rho}\right)^2 \times \left[1 + \frac{\rho}{R}k(R)\right]$$
$$g_y(R,\theta) = \left(\frac{R}{\rho}\right)^2 \times n \approx -\frac{R}{\rho} \times k(R)$$

Introduce the azimuthal dependence of k in the following way:

$$k(R,\theta) = \begin{cases} k_F(R), & \text{if } \theta \in \theta_F \\ k_D(R), & \text{if } \theta \in \theta_D \\ k_{drift}(R), & \text{if } \theta \in \theta_{drift} \end{cases}$$

• This yields:

$$\nu_x^2(R_E) = \sum_i \beta_i(R_E) - \sum_i \alpha_i(R_E) \times k_i(R_E)$$

$$\nu_y^2(R_E) = \sum_i \alpha_i(R_E) \times k_i(R_E) + \mathcal{F}^2 \left[1 + 2\tan^2(\xi) \right]$$

n

-1

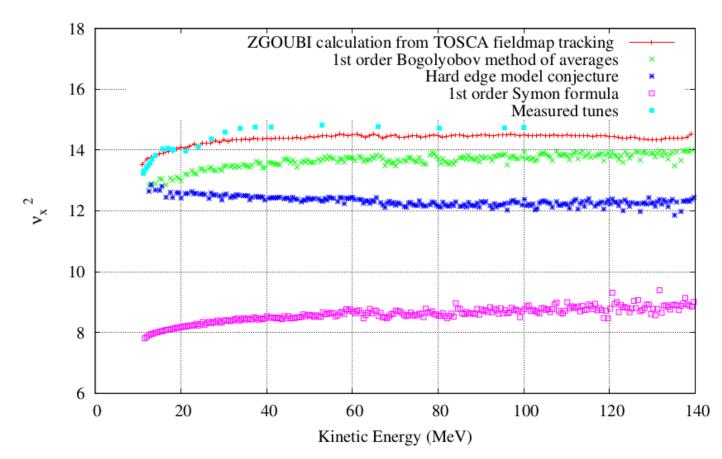
where

$$\alpha_i(R_E) = \frac{-1}{2\pi/N} \int_{\theta_i} \mu(R,\theta) d\theta = \frac{-1}{2\pi/N} \int_{\theta_i} \frac{R}{\rho} d\theta$$
$$\beta_i(R_E) = \frac{1}{2\pi/N} \int_{\theta_i} \mu(R,\theta)^2 d\theta = \frac{1}{2\pi/N} \int_{\theta_i} \left(\frac{R}{\rho}\right)^2 d\theta$$

Interpretation ...

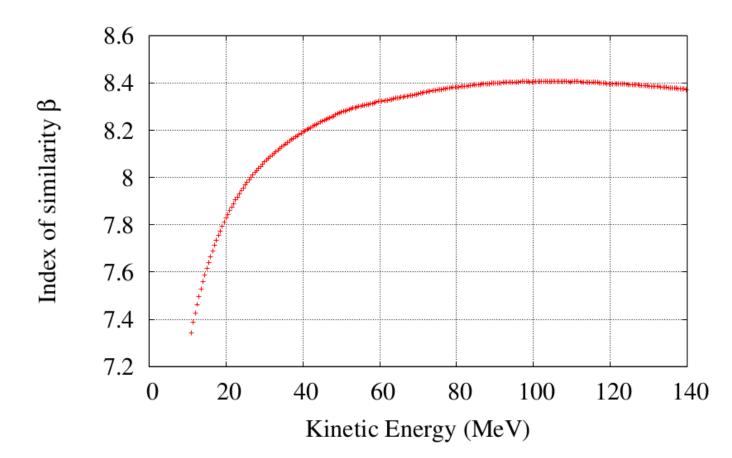
Application to the KURRI 150 MeV FFAG

 $\nu_x^2(E) = \beta_F(E) + 2\beta_D(E) + \beta_{drift}(E) - \alpha_F(E)k_F(E) - 2\alpha_D(E)k_D(E) - \alpha_{drift}(E)k_{drift}(E)$ $\nu_y^2(E) = \alpha_F(E)k_F(E) + 2\alpha_D(E)k_D(E) + \alpha_{drift}(E)k_{drift}(E)$



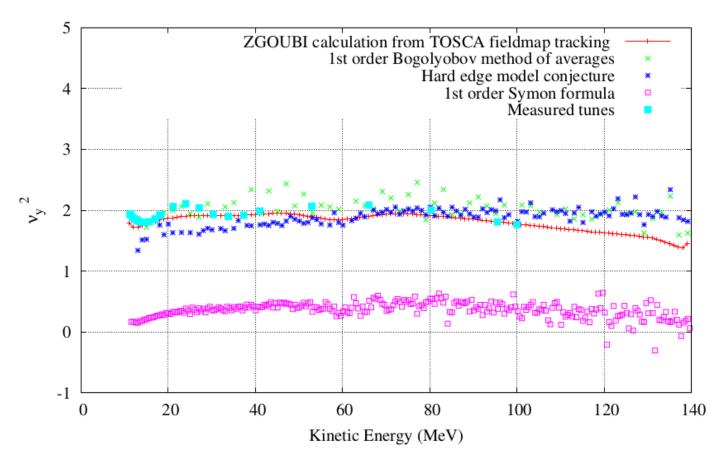
Index of similarity β

$$\beta_i(R_E) = \frac{1}{2\pi/N} \int_{\theta_i} \mu(R,\theta)^2 d\theta = \frac{1}{2\pi/N} \int_{\theta_i} \left(\frac{R}{\rho}\right)^2 d\theta$$

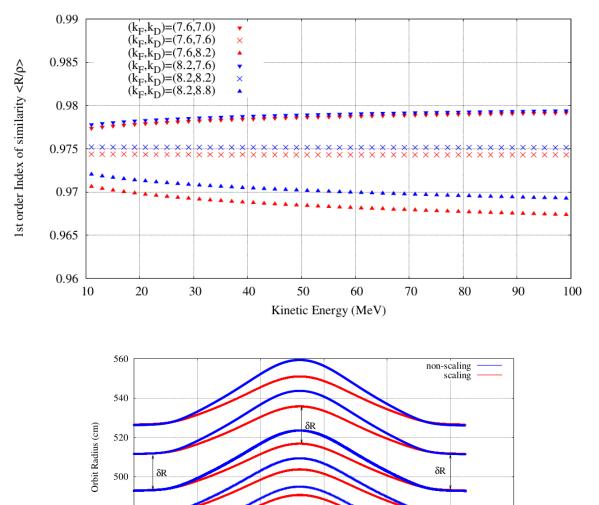


Application to the KURRI 150 MeV FFAG

 $\nu_x^2(E) = \beta_F(E) + 2\beta_D(E) + \beta_{drift}(E) - \alpha_F(E)k_F(E) - 2\alpha_D(E)k_D(E) - \alpha_{drift}(E)k_{drift}(E)$ $\nu_y^2(E) = \alpha_F(E)k_F(E) + 2\alpha_D(E)k_D(E) + \alpha_{drift}(E)k_{drift}(E)$



Scaling $(k_F = k_D)$ vs non-scaling $(k_F \neq k_D)$



480

460

0

0.1

0.2

0.3

Azimuthal angle θ (rad)

0.4

0.5

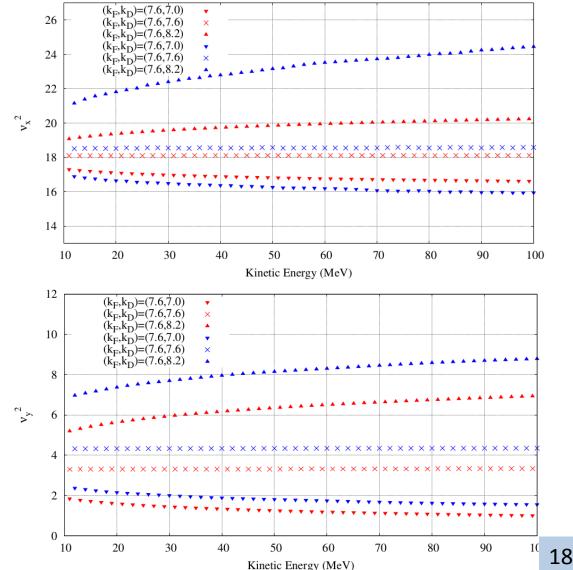
0.6

The F and D magnets show antagonistic behavior

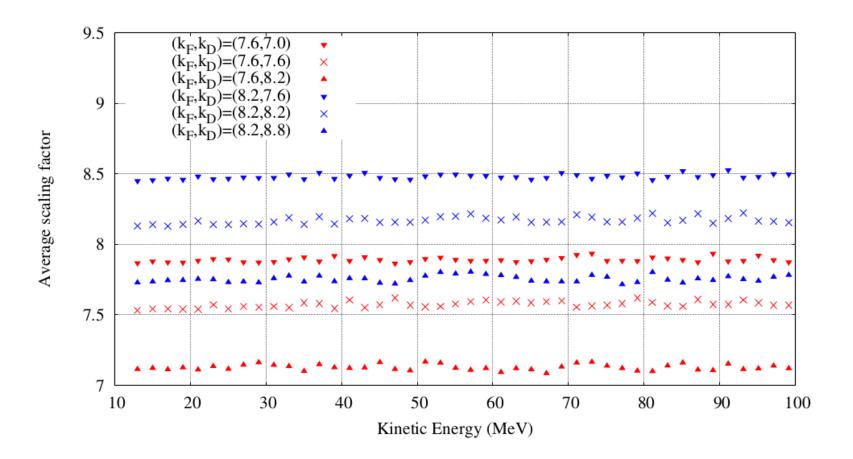
Comparison of the Zgoubi results (blue) with the 1st order bogoliubov method of averages show good agreement.

Dpending on whether $k_F > k_D$ or the opposite, the tune exhibits antagonistic behavior.

Thus, in presence of systematic errors of the field, i.e $k_D > k_F$ for instance, the idea would be to introduce the opposite error by implementing trim coils every two sectors. The superperiod becomes DFD-DFD



Average scaling factor



Having the average scaling factor k constant is not a sufficient condition to obtain a fixed tune machine

Thank you